



Universität Hamburg
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Course 2: Basic Technologies

Part VI: Synchrotron Sources / FEL

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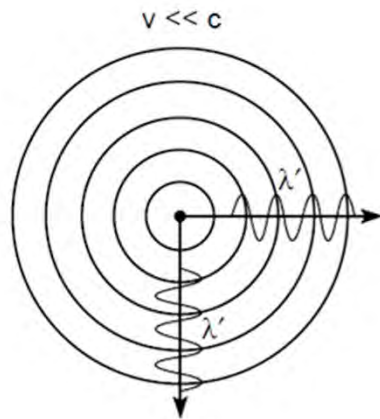
Synchrotron Radiation



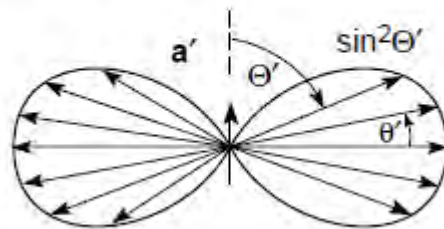
Facts and figures

- Ring accelerator for electrons and positrons
- Length: 2304 metres
- Commissioning: 1978
- 1978-1986: particle physics
- 1987-2007: pre-accelerator for HERA and X-ray radiation source
- Since 2009: most brilliant storage-ring-based X-ray source in the world
- Start of user operation: 2009
- 14 experimental stations with up to 30 instruments

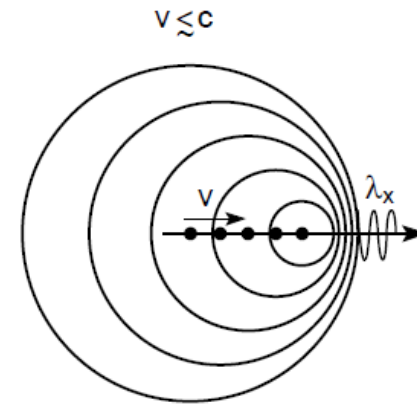
relativistic electrons



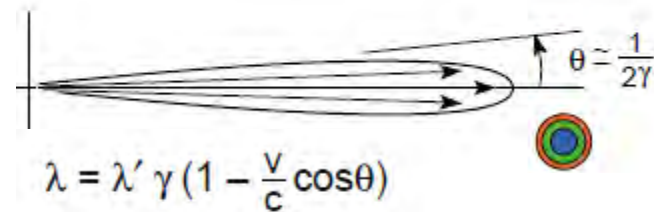
Frame moving with electron



$$\lambda = \lambda' \left(1 - \frac{v}{c} \cos \theta\right)$$



Laboratory frame of reference



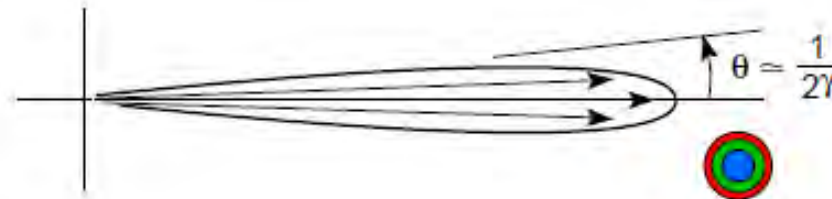
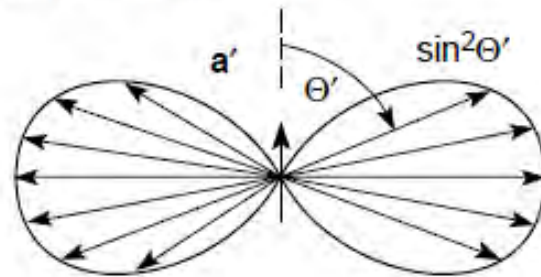
$$\lambda = \lambda' \gamma \left(1 - \frac{v}{c} \cos \theta\right)$$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

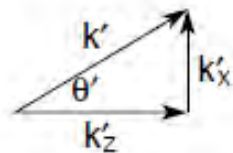
$$\tan \theta = \frac{\sin \theta'}{\gamma (\beta + \cos \theta')}$$

Relativistic Electrons Radiate in a Narrow Forward Cone

Dipole radiation

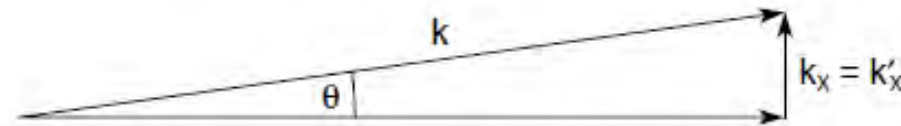


Frame of reference moving with electrons



$$k' = 2\pi/\lambda'$$

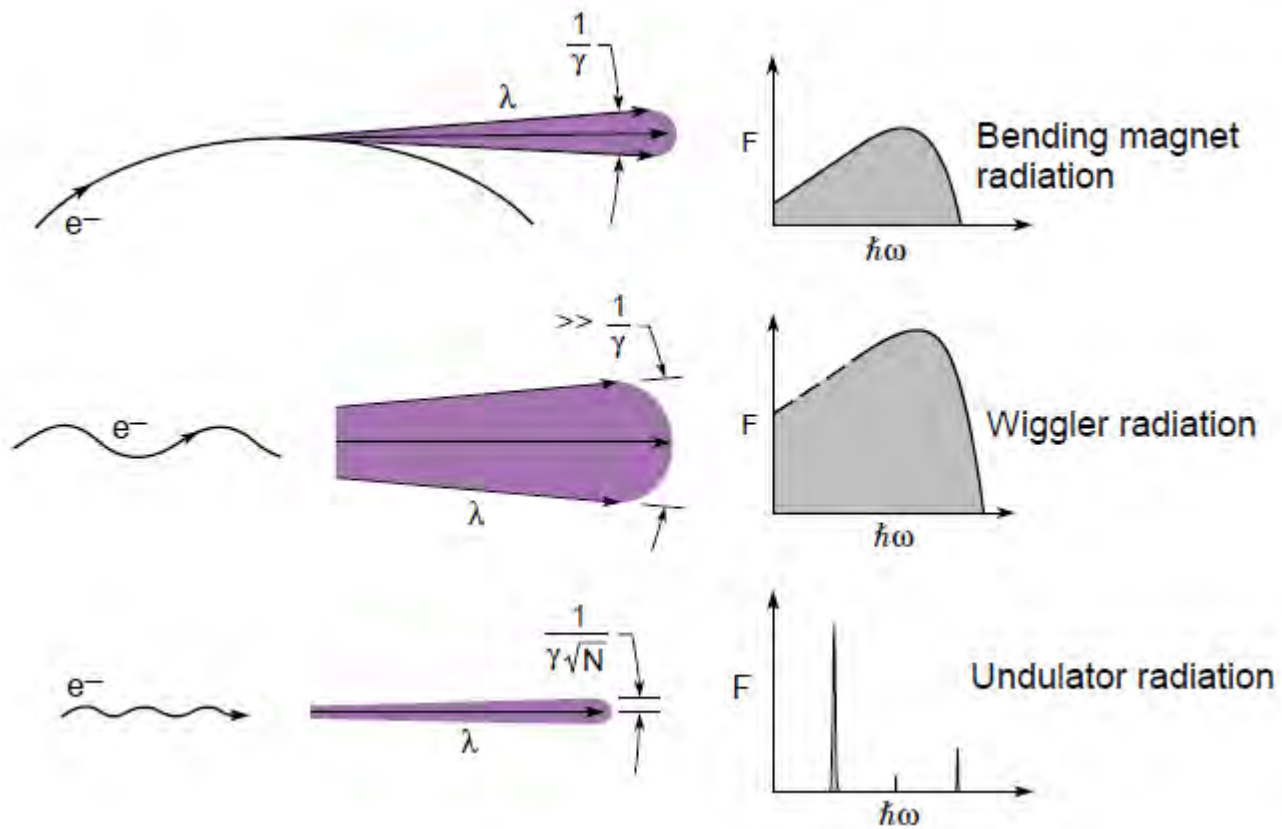
Lorentz transformation



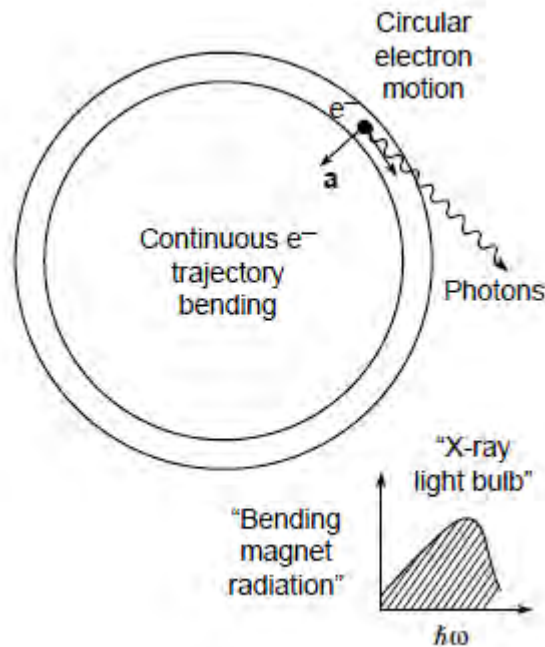
$$k_z = 2\gamma k'_z \text{ (Relativistic Doppler shift)}$$

$$\theta = \frac{k_x}{k_z} = \frac{k'_x}{2\gamma k'_z} = \frac{\tan\theta'}{2\gamma} = \frac{1}{2\gamma}$$

Three types of radiation



1. Bending magnet radiation (... „old style“)



$$R = \frac{\gamma m v}{e B} \simeq \frac{\gamma m c}{e B}$$

The Lorentz force for a relativistic electron in a constant magnetic field is

$$\mathbf{F} = \frac{d\mathbf{p}}{dt} = -e\mathbf{v} \times \mathbf{B}$$

where $\mathbf{p} = \gamma m \mathbf{v}$. In a fixed magnetic field the rate of change of electron energy is

$$\frac{dE_e}{dt} = \mathbf{v} \cdot \mathbf{F} = \underbrace{-e\mathbf{v} \cdot (\mathbf{v} \times \mathbf{B})}_{=0}$$

thus with $E_e = \gamma m c^2$

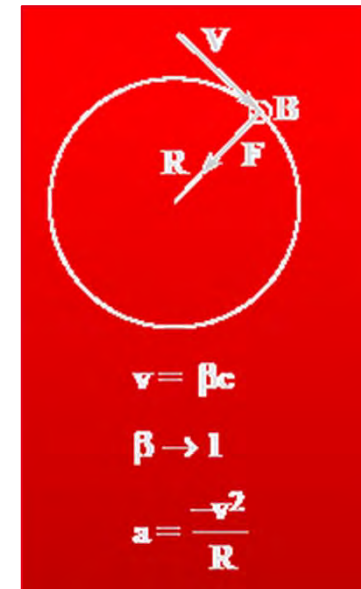
$$\frac{dE_e}{dt} = \frac{d}{dt}(\gamma m c^2) = 0$$

$\therefore \gamma = \text{constant}$

and the force equation becomes

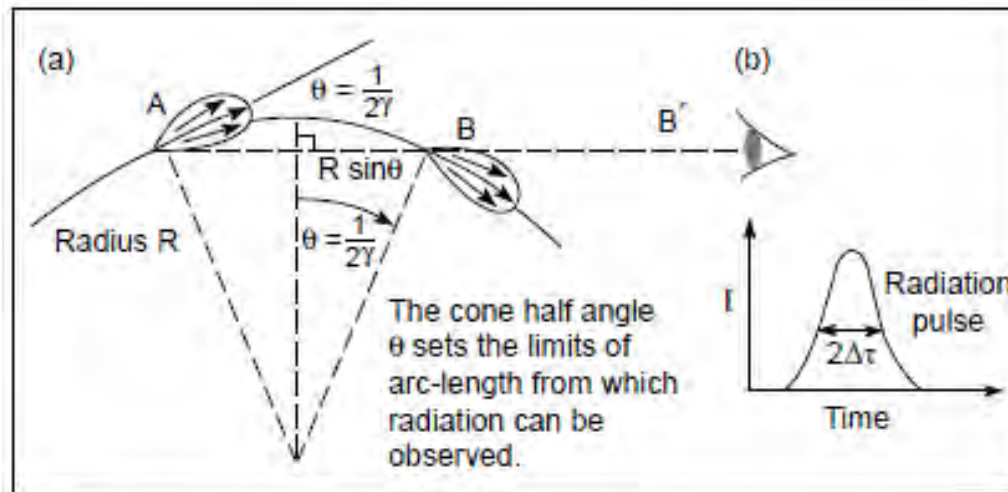
$$\frac{d\mathbf{p}}{dt} = \gamma m \frac{d\mathbf{v}}{dt} = -e\mathbf{v} \times \mathbf{B}$$

$$\gamma m \left(-\frac{v^2}{R} \right) = -e v B$$



http://henke.lbl.gov/optical_constants/bend2.html

Bending Magnet Radiation



$$2 \Delta \tau = \tau_e - \tau_r$$

$$2 \Delta \tau = \frac{\text{arc length}}{v} - \frac{\text{radiation path}}{c}$$

$$2 \Delta \tau \simeq \frac{R \cdot 2\theta}{v} - \frac{2R \sin \theta}{c}$$

With $\theta \simeq 1/2\gamma$, $\sin \theta \simeq \theta$

$$2 \Delta \tau \simeq \frac{R}{\gamma v} - \frac{R}{\gamma c} = \frac{R}{\gamma} \left(\frac{1}{v} - \frac{1}{c} \right)$$

With $v = \beta c$

$$2 \Delta \tau \simeq \frac{R}{\gamma \beta c} (1 - \beta) \quad \text{but} \quad (1 - \beta) \simeq \frac{1}{2\gamma^2} \quad \text{and} \quad R \simeq \frac{\gamma m c}{e B}$$

$$2 \Delta \tau = \frac{m}{2eB\gamma^2}$$

Bending magnet radiation revisited

From Heisenberg's Uncertainty Principle for rms pulse duration and photon energy

$$\Delta E \cdot \Delta \tau \geq \hbar/2$$

thus

$$\Delta E \geq \frac{\hbar}{2\Delta \tau}$$

$$\Delta E \geq \frac{\hbar}{m/2eB\gamma^2}$$

Thus the single-sided rms photon energy width (uncertainty) is

$$\Delta E \geq \frac{2e\hbar B\gamma^2}{m}$$

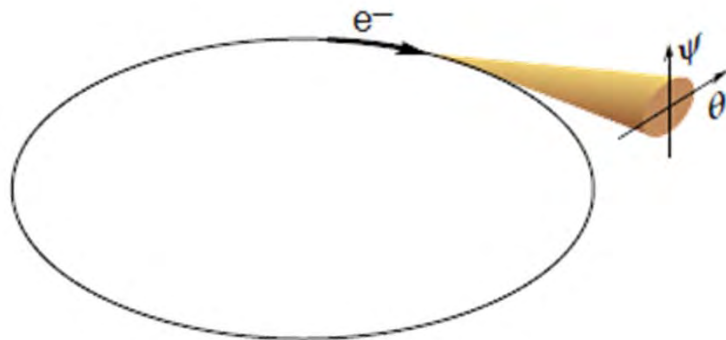
A more detailed description of bending magnet radius finds the critical photon energy

$$E_c = \hbar\omega_c = \frac{3e\hbar B\gamma^2}{2m}$$

In practical units the critical photon energy is

$$E_c(\text{keV}) = 0.6650 E_e^2(\text{GeV}) B(\text{T})$$

bending magnet - summary

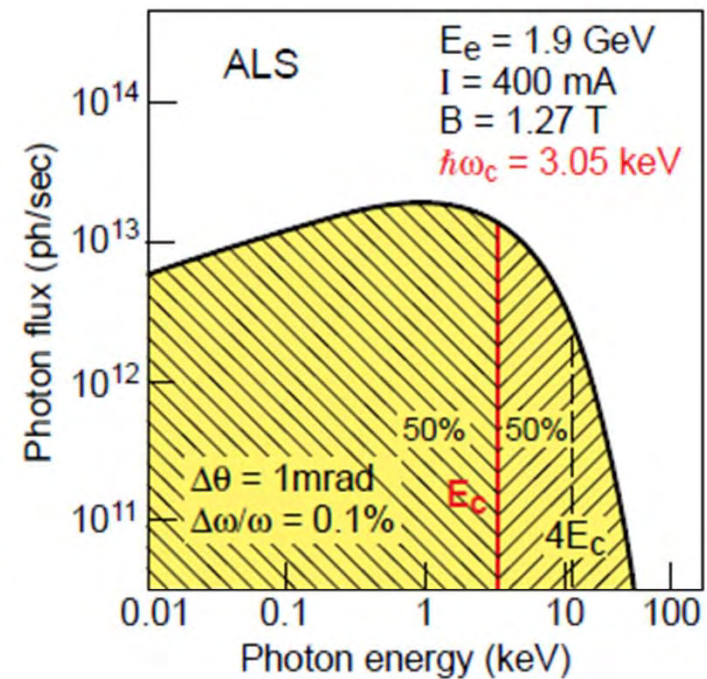


$$E_c = \hbar\omega_c = \frac{3e\hbar B\gamma^2}{2m} \quad (5.7a)$$

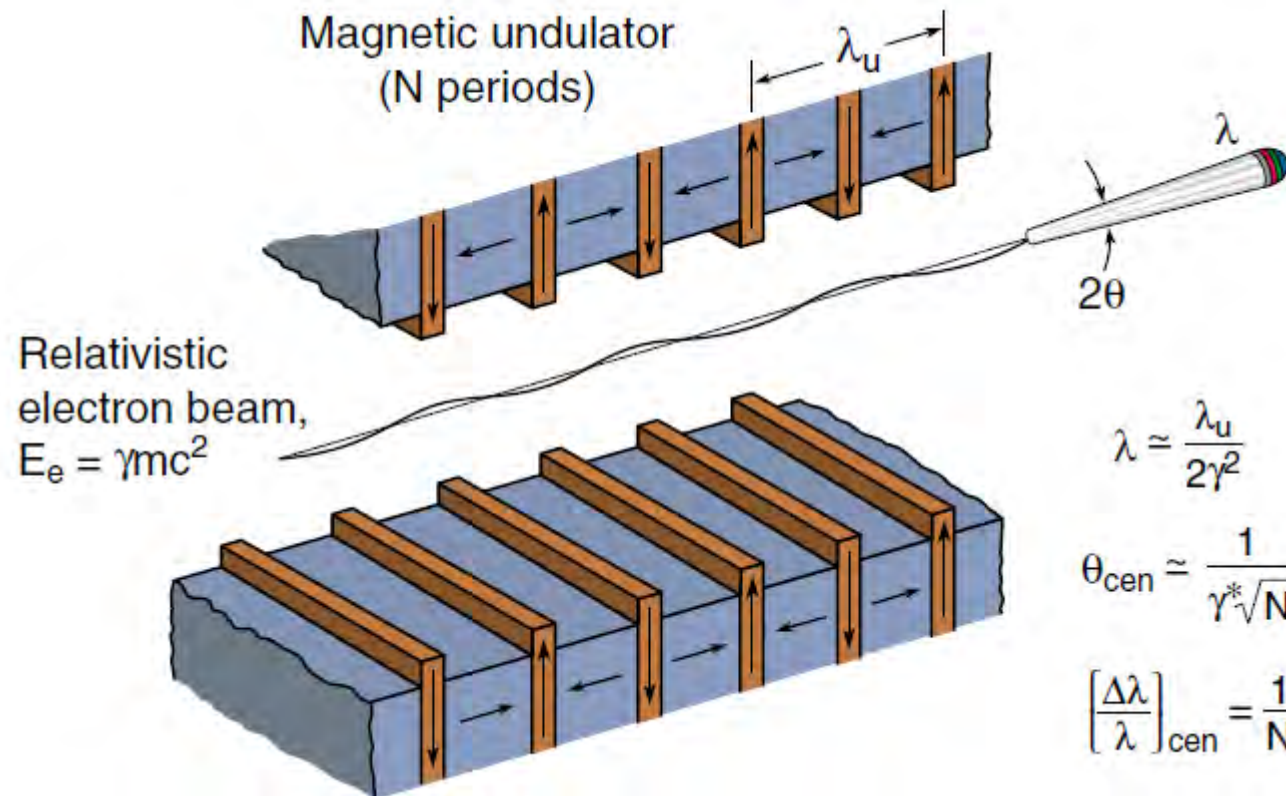
$$E_c(\text{keV}) = 0.6650 E_e^2(\text{GeV}) B(\text{T}) \quad (5.7b)$$

$$\frac{d^2 F_B}{d\theta d\omega/\omega} = 2.46 \times 10^{13} E_e(\text{GeV}) I(\text{A}) G_1(E/E_c) \frac{\text{photons/s}}{\text{mrad} \cdot (0.1\% \text{BW})} \quad (5.8)$$

- | | |
|-----------------------|---|
| <u>Advantages:</u> | <ul style="list-style-type: none"> • covers broad spectral range • least expensive • most accessible |
| <u>Disadvantages:</u> | <ul style="list-style-type: none"> • limited coverage of hard x-rays • not as bright as undulator |

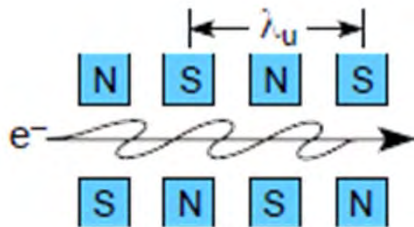


Undulator



undulator radiation

Laboratory Frame of Reference

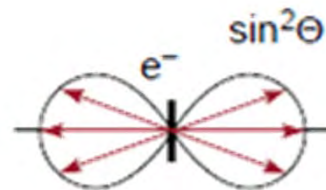


$$E = \gamma mc^2$$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$N = \# \text{ periods}$

Frame of moving e-



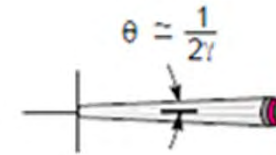
e^- radiates at the Lorentz contracted wavelength:

$$\lambda' = \frac{\lambda_u}{\gamma}$$

Bandwidth:

$$\frac{\lambda'}{\Delta\lambda'} \approx N$$

Frame of observer



Doppler shortened wavelength on axis:

$$\lambda = \lambda' \gamma (1 - \beta \cos \theta)$$

$$\lambda = \frac{\lambda_u}{2\gamma^2} (1 + \gamma^2 \theta^2)$$

Accounting for transverse motion due to the periodic magnetic field:

$$\lambda = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + \gamma^2 \theta^2 \right)$$

where $K = eB_0\lambda_u / 2\pi mc$

where does the $\lambda = \lambda_u/2\gamma^2$ come from?

The electron “sees” a Lorentz contracted period

$$\lambda' = \frac{\lambda_u}{\gamma}$$

and emits radiation in its frame of reference at frequency

$$f' = \frac{c}{\lambda'} = \frac{c\gamma}{\lambda_u}$$

Observed in the laboratory frame of reference, this radiation is Doppler shifted to a frequency

$$f = \frac{f'}{\gamma(1 - \beta \cos \theta)} = \frac{c}{\lambda_u(1 - \beta \cos \theta)}$$

On-axis ($\theta = 0$) the observed frequency is

$$f = \frac{c}{\lambda_u(1 - \beta)}$$

$$f = \frac{c}{\lambda_u(1 - \beta)}$$

By definition $\gamma = \frac{1}{\sqrt{1 - \beta^2}}$; $\gamma^2 = \frac{1}{(1 - \beta)(1 + \beta)} \approx \frac{1}{2(1 - \beta)}$

thus

$$f = \frac{2\gamma^2 c}{\lambda_u}$$

and the observed wavelength is

$$\lambda = \frac{c}{f} = \frac{\lambda_u}{2\gamma^2}$$

With electrons executing N oscillations as they traverse the periodic magnet structure, and thus radiating a wavetrain of N cycles, it is of interest to know what angular cone contains radiation of relative spectral bandwidth

$$\frac{\Delta\lambda}{\lambda} = \frac{1}{N} \quad (5.14)$$

Write the undulator equation twice, once for on-axis radiation ($\theta = 0$) and once for length-shifted radiation off-axis at angle θ :

$$\lambda_0 + \Delta\lambda = \frac{\lambda_u}{2\gamma^2} (1 + \gamma^2\theta^2)$$

$$\lambda_0 = \frac{\lambda_u}{2\gamma^2}$$

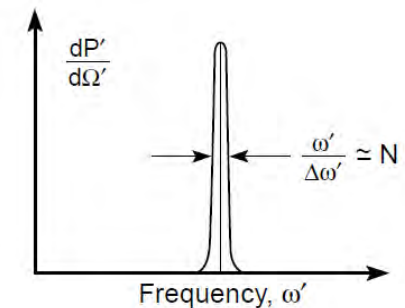
divide and simplify to

$$\frac{\Delta\lambda}{\lambda} \simeq \gamma^2\theta^2 \quad (5.13)$$

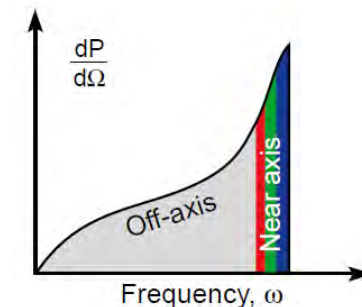
Combining the two equations (5.13 and 5.14)

defines θ_{cen} : $\gamma^2\theta_{\text{cen}}^2 \equiv \frac{1}{N}$, which gives $\theta_{\text{cen}} \simeq \frac{1}{\gamma\sqrt{N}}$ (5.15)

This is the half-angle of the “central radiation cone”, defined as containing radiation of $\Delta\lambda/\lambda = 1/N$.



Execution of N electron oscillations produces a transform-limited spectral bandwidth, $\Delta\omega'/\omega' = 1/N$.



The Doppler frequency shift has a strong angle dependence, leading to lower photon energies off-axis.

radiation power of an undulator

$$\lambda_x = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + \gamma^2 \theta^2 \right)$$

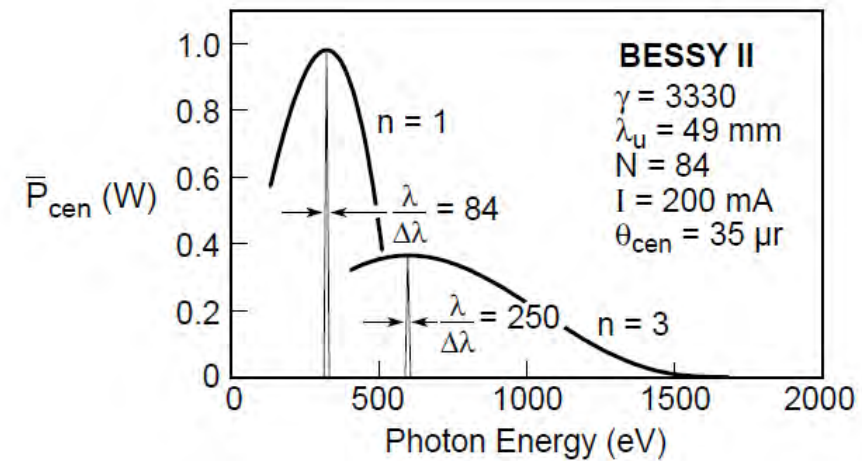
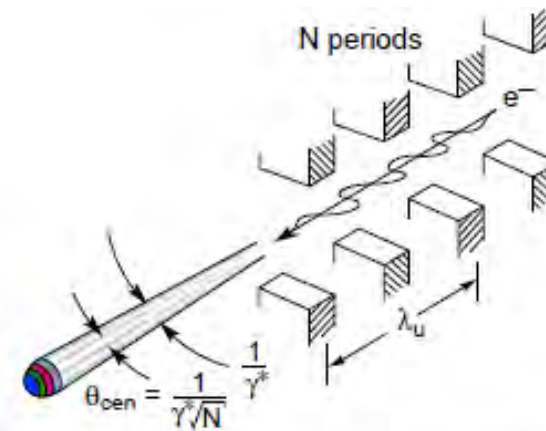
$$\bar{P}_{\text{cen}} = \frac{\pi e \gamma^2 I}{\epsilon_0 \lambda_u} \frac{K^2}{\left(1 + \frac{K^2}{2} \right)^2} f(K)$$

$$\theta_{\text{cen}} = \frac{1}{\gamma^* \sqrt{N}}$$

$$\left(\frac{\Delta\lambda}{\lambda} \right)_{\text{cen}} = \frac{1}{N}$$

$$K = \frac{e B_0 \lambda_u}{2\pi m_0 c}$$

$$\gamma^* = \gamma / \sqrt{1 + \frac{K^2}{2}}$$



Magnetic fields in the periodic undulator cause the electrons to oscillate and thus radiate. These magnetic fields also slow the electrons axial (z) velocity somewhat, reducing both the Lorentz contraction and the Doppler shift, so that the observed radiation wavelength is not quite so short. The force equation for an electron is

$$\frac{d\mathbf{p}}{dt} = -e(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

where $\mathbf{p} = \gamma m \mathbf{v}$ is the momentum. The radiated fields are relatively weak so that

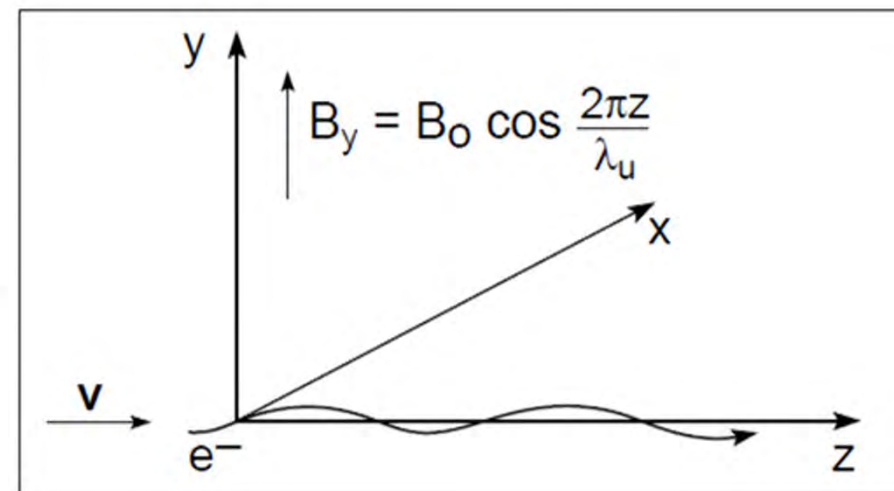
$$\frac{d\mathbf{p}}{dt} \approx -e(\mathbf{v} \times \mathbf{B})$$

Taking to first order $v \approx v_z$, motion in the x -direction is

$$m\gamma \frac{dv_x}{dt} = +ev_z B_y$$

$$m\gamma \frac{dv_x}{dt} = e \frac{dz}{dt} \cdot B_0 \cos\left(\frac{2\pi z}{\lambda_u}\right) \quad (0 \leq z \leq N\lambda_u)$$

$$m\gamma dv_x = e dz B_0 \cos\left(\frac{2\pi z}{\lambda_u}\right)$$



$$m\gamma dv_x = e dz B_0 \cos\left(\frac{2\pi z}{\lambda_u}\right)$$

integrating both sides

$$m\gamma v_x = e B_0 \frac{\lambda_u}{2\pi} \int \cos\left(\frac{2\pi z}{\lambda_u}\right) \cdot d\left(\frac{2\pi z}{\lambda_u}\right)$$

$$m\gamma v_x = \frac{e B_0 \lambda_u}{2\pi} \sin\left(\frac{2\pi z}{\lambda_u}\right)$$

$$v_x = \frac{Kc}{\gamma} \sin\left(\frac{2\pi z}{\lambda_u}\right)$$

$$K \equiv \frac{e B_0 \lambda_u}{2\pi m c} = 0.9337 B_0(\text{T}) \lambda_u(\text{cm})$$

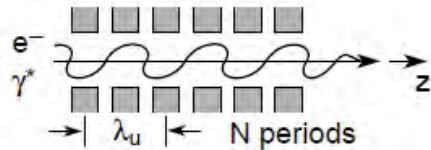
is the non-dimensional “magnetic deflection parameter.”

The “deflection angle”, θ , is

$$\theta = \frac{v_x}{v_z} \approx \frac{v_x}{c} = \frac{K}{\gamma} \sin k_u z$$

calculating the radiation power of the central cone

x, z, t laboratory frame of reference



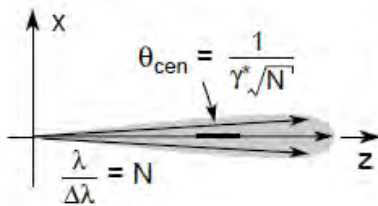
Determine x, z, t motion:

$$\frac{d\mathbf{p}}{dt} = -e(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

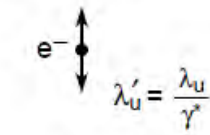
$$m\gamma \frac{dv_x}{dt} = e \frac{dz}{dt} B_0 \cos \frac{2\pi z}{\lambda_u}$$

$$v_x(t); a_x(t) = \dots$$

$$v_z(t); a_z(t) = \dots$$



x', z', t' frame of reference moving with the average velocity of the electron
 Lorentz transformation

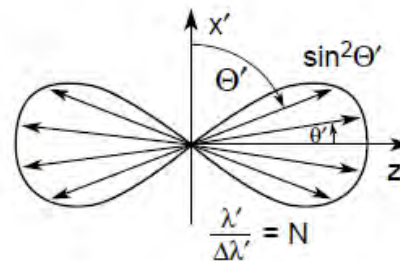


x', z', t' motion
 $a'(t')$ acceleration

Dipole radiation:

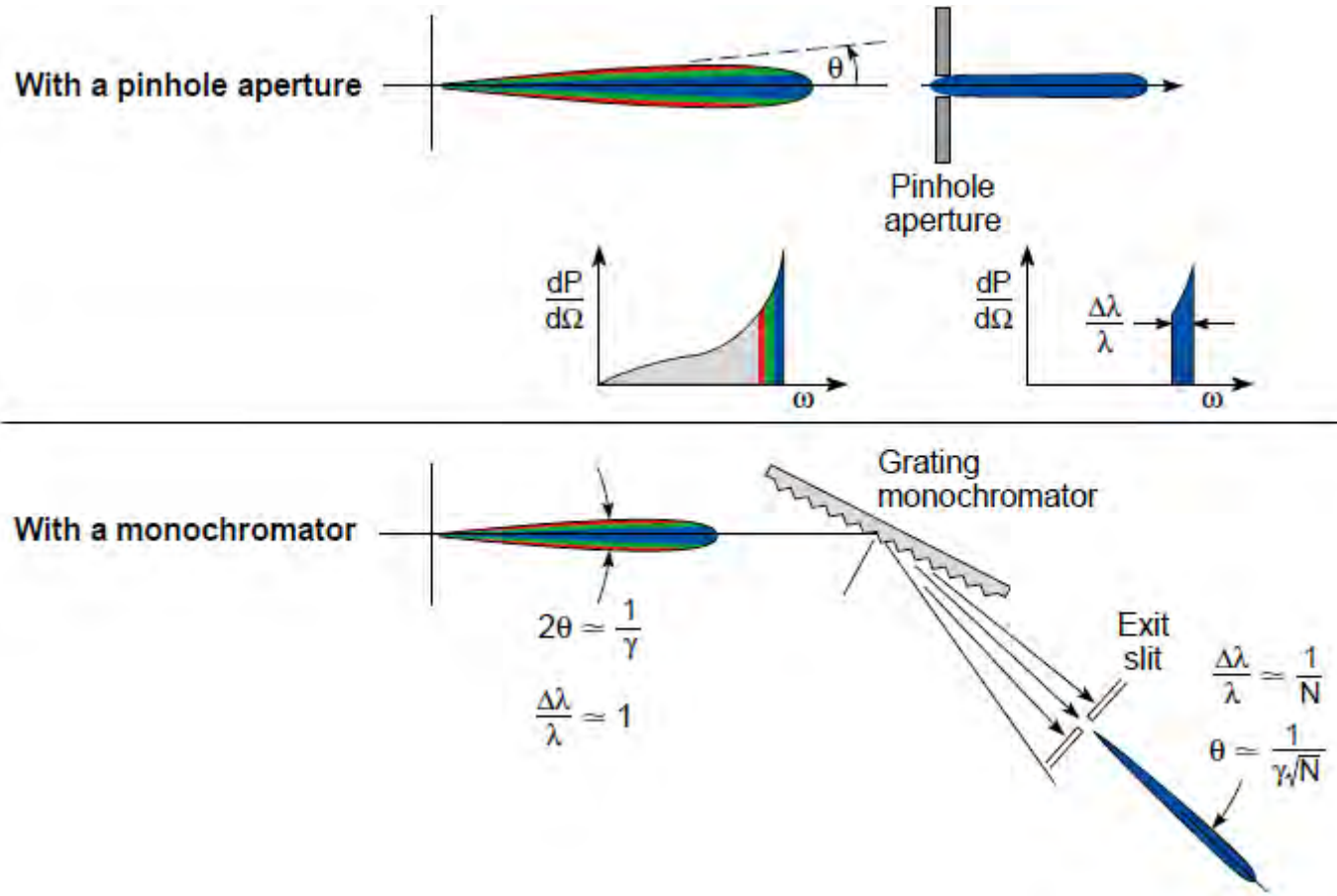
$$\frac{dP'}{d\Omega'} = \frac{e^2 a'^2 \sin^2 \Theta'}{16\pi^2 \epsilon_0 c^3}$$

$$\frac{dP'}{d\Omega'} = \frac{e^2 c \gamma^2}{4\epsilon_0 \lambda_u^2} \frac{K^2}{(1 + K^2/2)^2} (1 - \sin^2 \theta' \cos^2 \phi') \cos^2 \omega'_u t'$$



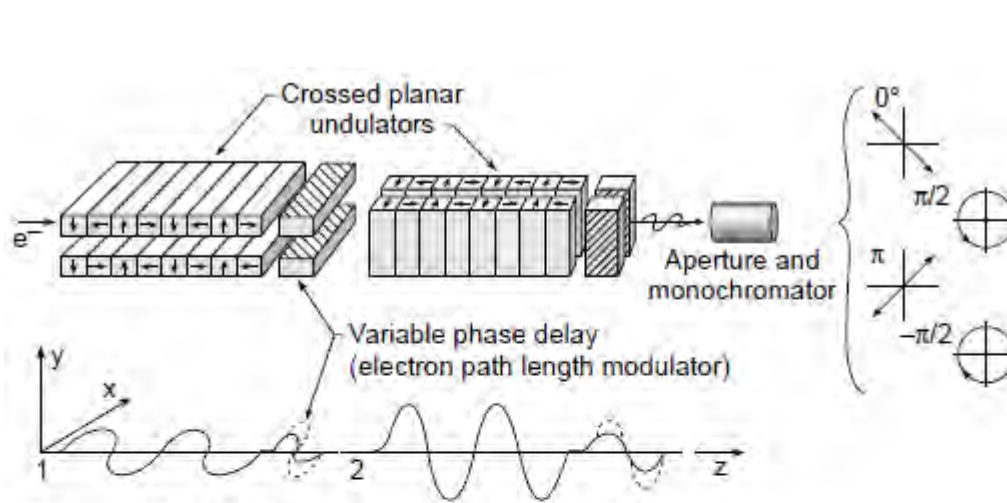
$$\bar{P}_{\text{cen}} = \frac{\pi e \gamma^2 I}{\epsilon_0 \lambda_u} \frac{K^2}{(1 + \frac{K^2}{2})^2} f(K)$$

bandwith of the undulator radiation

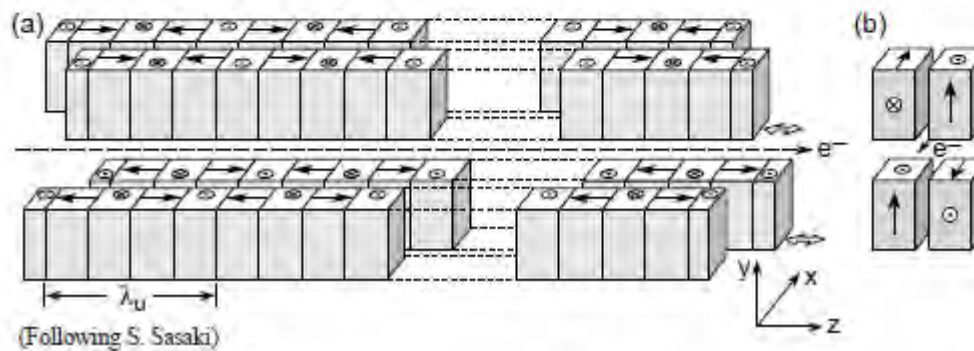
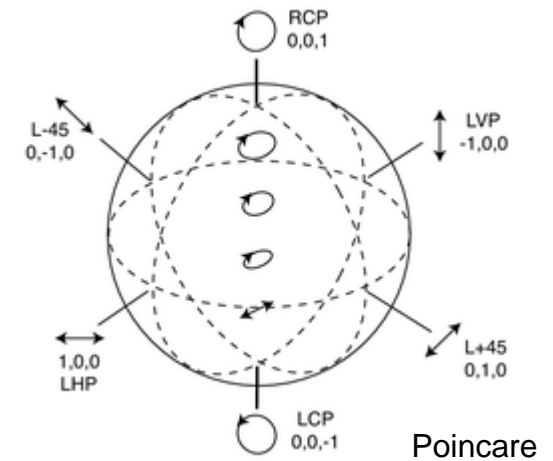


The Narrow (1/N) Spectral Bandwidth of Undulator Radiation Can be Recovered in Two Ways

Polarisation properties

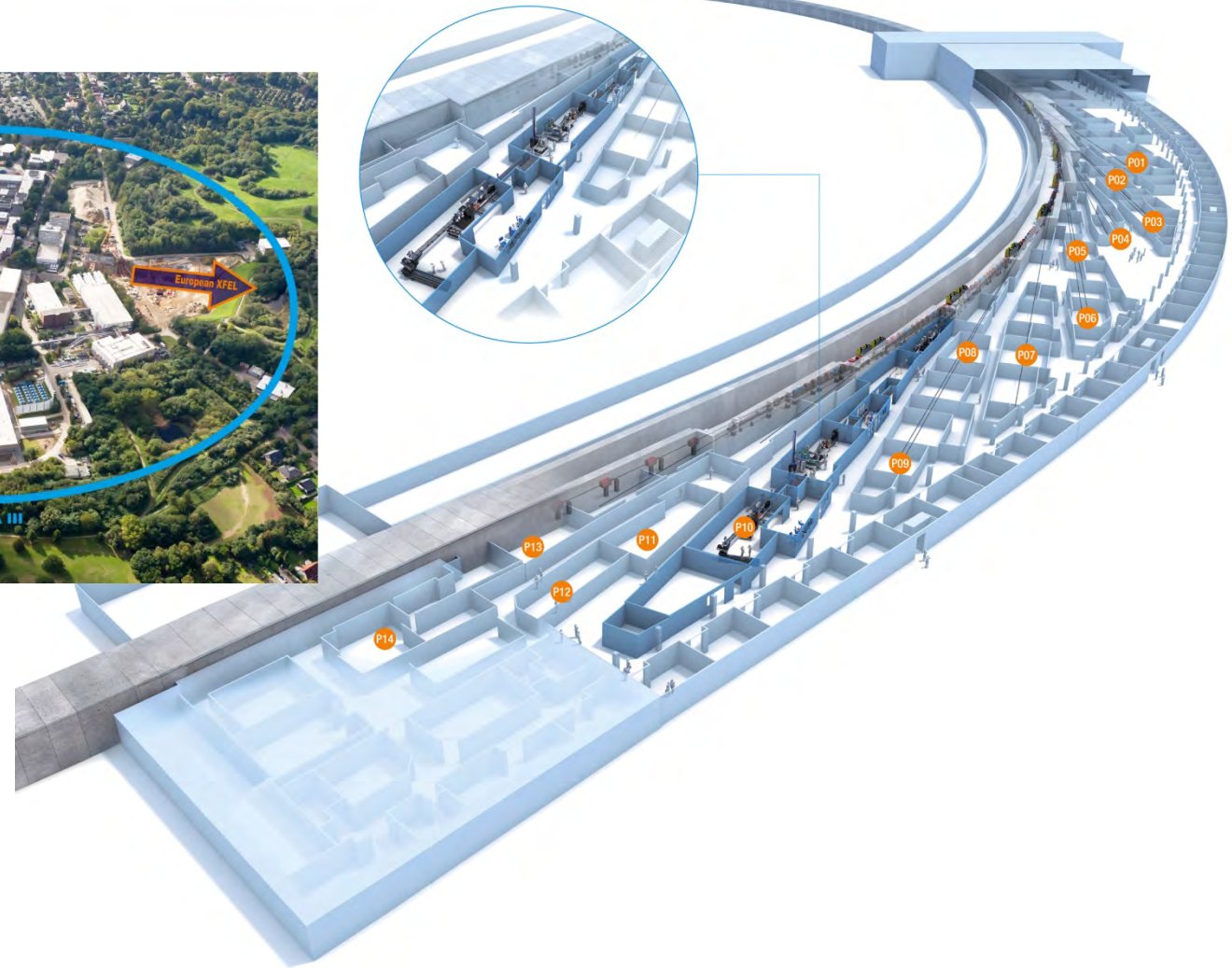
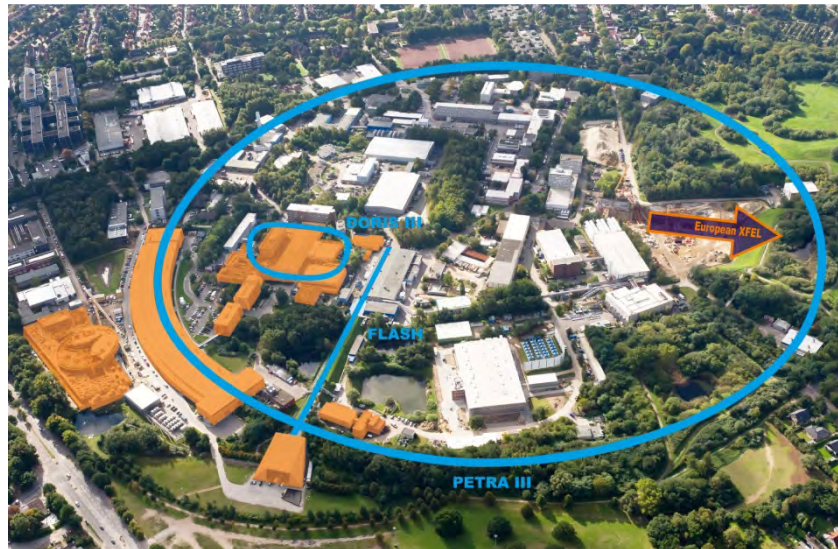


(Courtesy of Kwang-Je Kim)



(Following S. Sasaki)

PETRA III



Benefits of synchrotron radiation

Spectrum (tunability, bandwidth...)

Prediction

Polarisation (very clean)

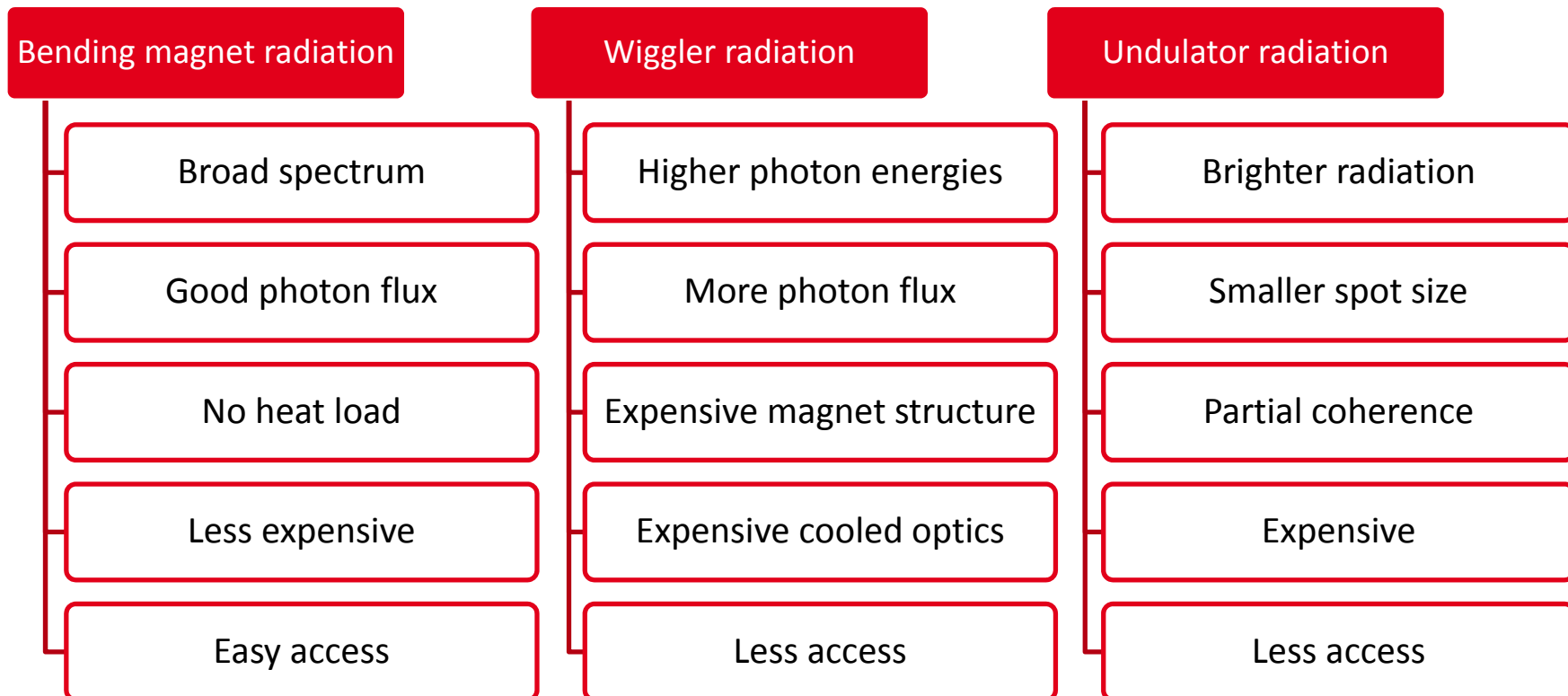
time structure (electron bunches)

UHV compatible

Intensity

High optical power with very good focusability

Comparison



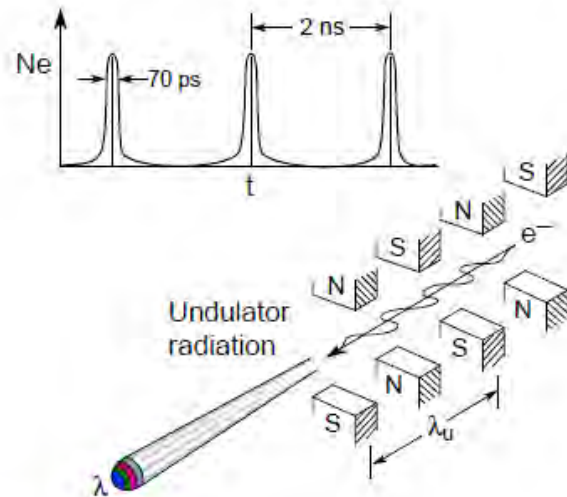
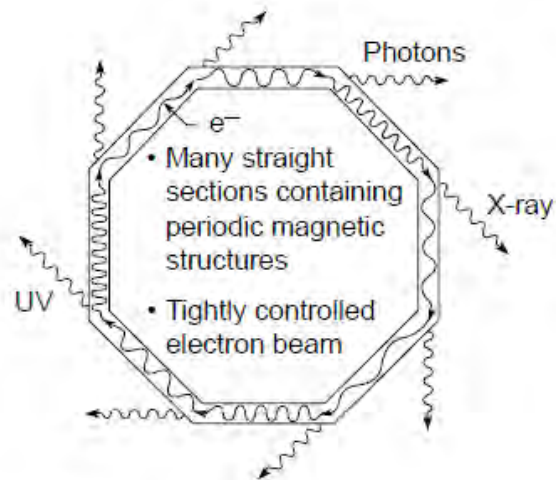
Further information

Prof. David Attwood

Soft X-Ray and Extreme Ultraviolet Radiation,
Cambridge University Press, 1999

<http://ast.coe.berkeley.edu/srms/>

Synchrotron radiation sources



Bending Magnet:

$$\hbar\omega_c = \frac{3e\hbar B\gamma^2}{2m}$$

Wiggler:

$$\hbar\omega_c = \frac{3e\hbar B\gamma^2}{2m}$$

$$n_c = \frac{3K}{4} \left(1 + \frac{K^2}{2}\right)$$

$$P_T = \frac{\pi e K^2 \gamma^2 I N}{3\epsilon_0 \lambda_u}$$

Undulator:

$$\lambda = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + \gamma^2\theta^2\right)$$

$$K = \frac{eB_0\lambda_u}{2\pi mc}$$

$$\theta_{cen} = \frac{1}{\gamma^* \sqrt{N}}$$

$$\left. \frac{\Delta\lambda}{\lambda} \right|_{cen} = \frac{1}{N}$$

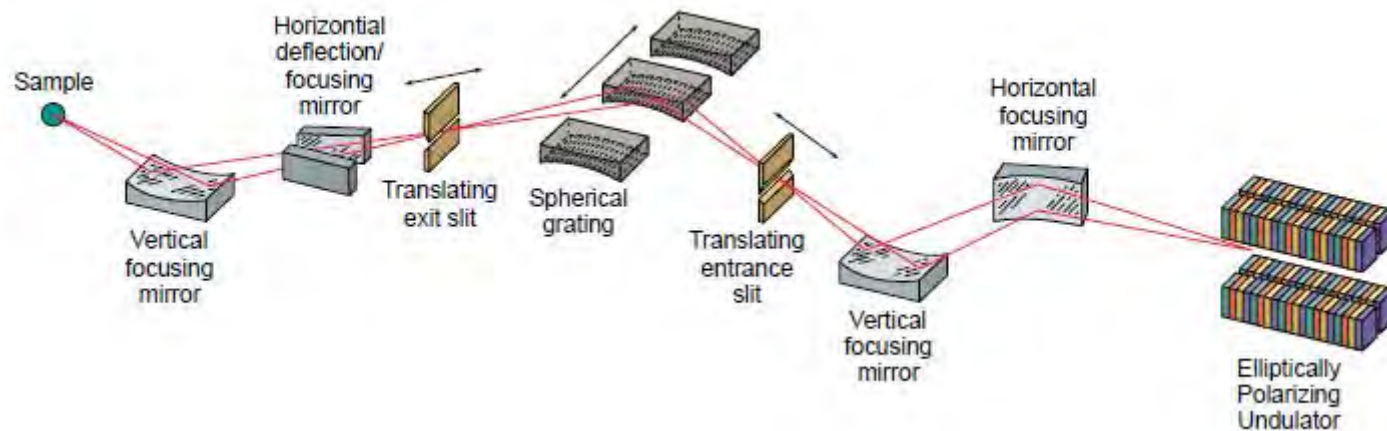
$$\bar{P}_{cen} = \frac{\pi e \gamma^2 I}{\epsilon_0 \lambda_u} \frac{K^2}{\left(1 + \frac{K^2}{2}\right)^2} f(K)$$

For a particle moving at a velocity v with a total energy E

$$E = \gamma m_0 c^2$$

$$\text{with } \gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

short reminder...





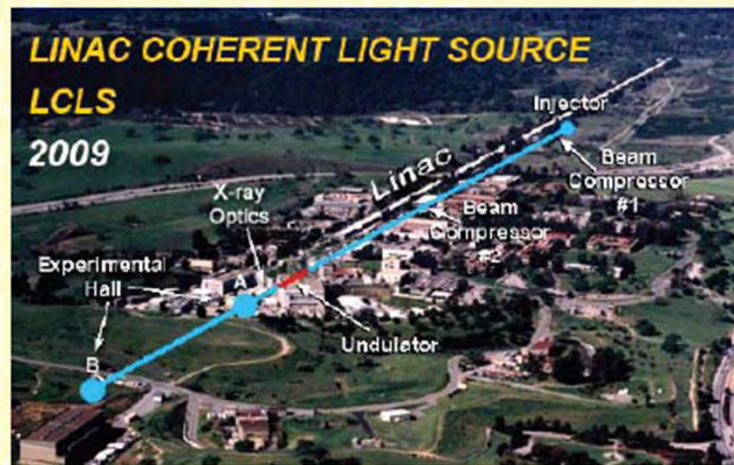
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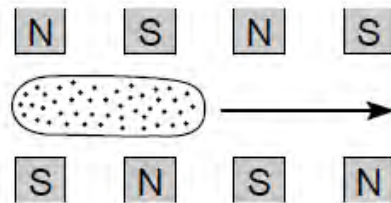
FEL – Free Electron Lasers

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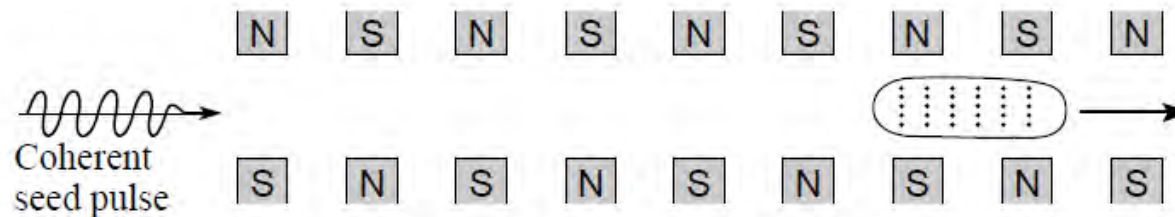
Some Free Electron Lasers



Undulator vs. FEL



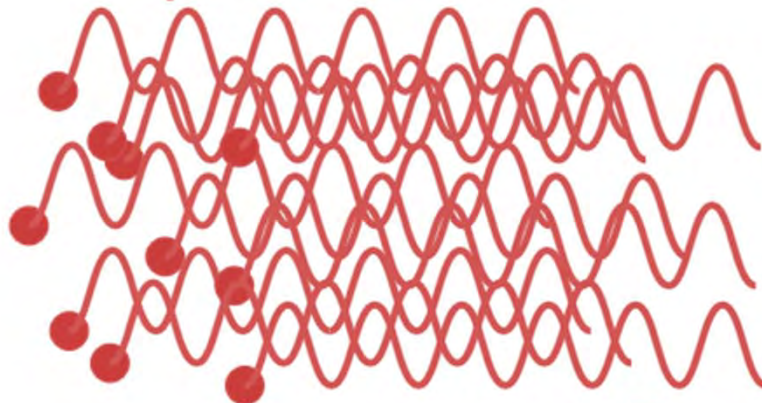
Undulator – uncorrelated electron positions, radiated fields uncorrelated, intensities add, limited coherence, power $\sim N$.



Free Electron Laser (FEL) – very long undulator, electrons are “microbunched” by their own radiated fields into strongly correlated waves of electrons, all radiated electric fields now add, spatially coherent, power $\sim N^2$

SR or ERL

Spontaneous Radiation



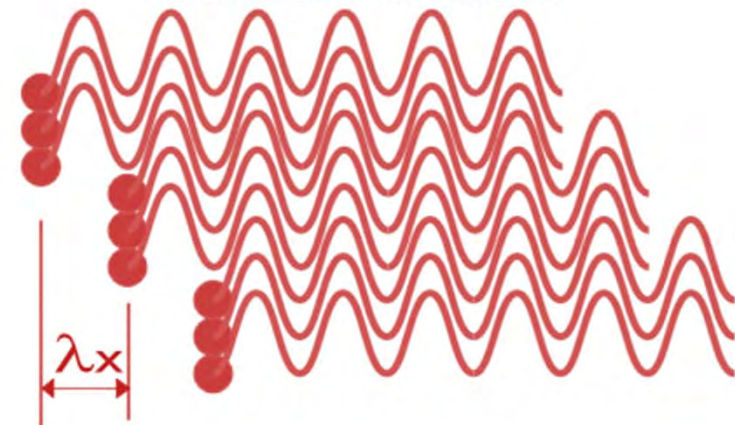
N -electrons
random distribution

$$E_{spt} \sim \sqrt{N} E_1$$

$$P_{spt} \sim N P_1$$

FEL: Free Electron Laser

Coherent Radiation

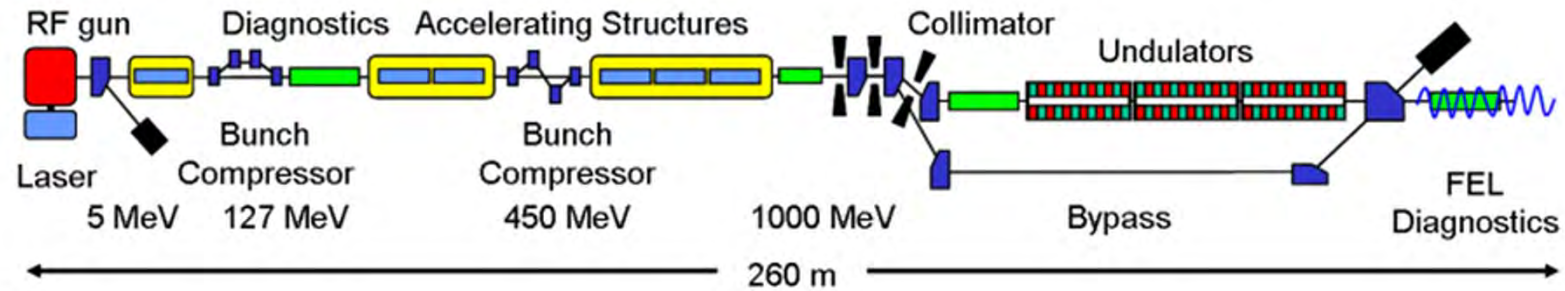


N -electrons
micro-bunched

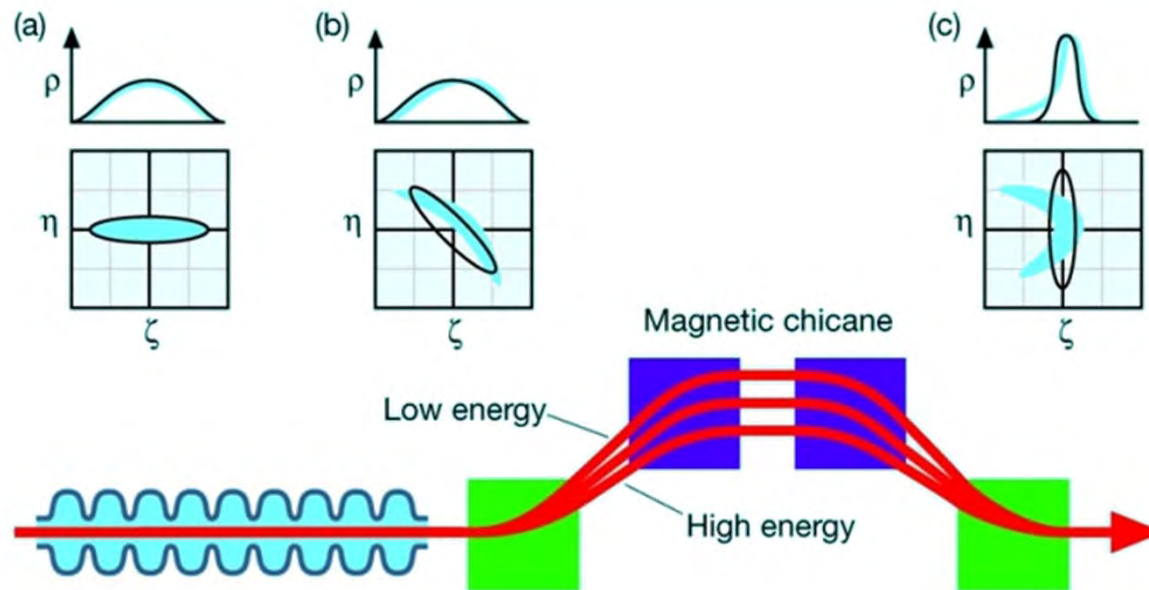
$$E_{coherent} \sim N E_1$$

$$P_{coherent} \sim N^2 P_1$$

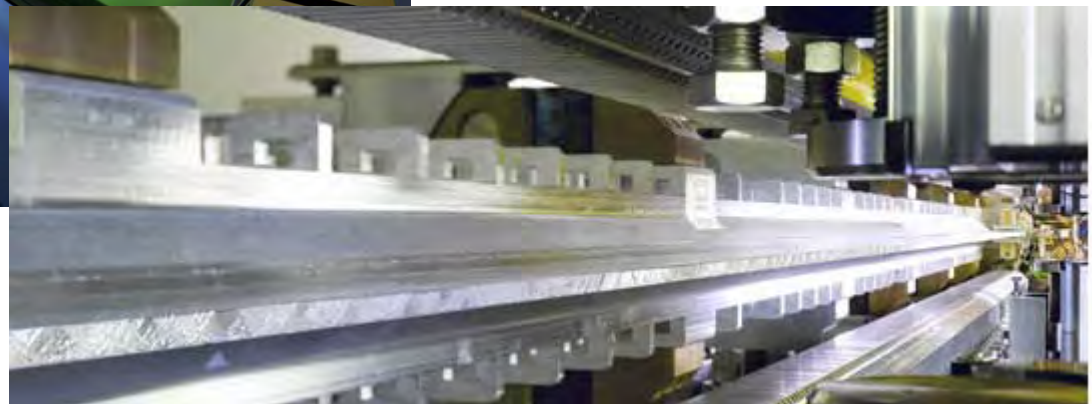
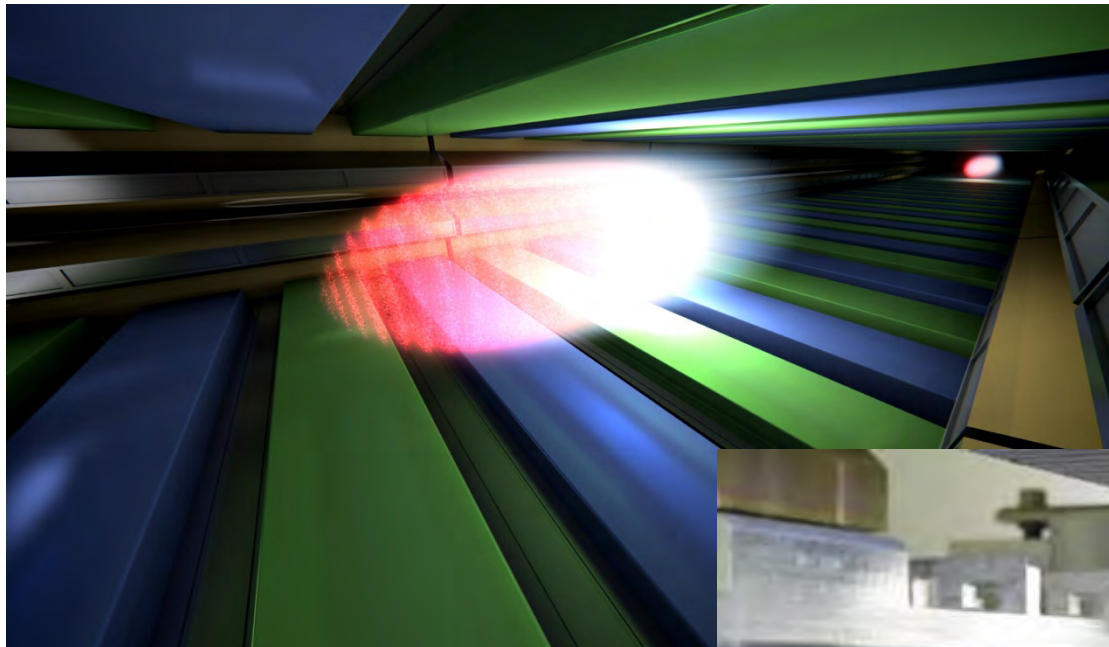
Flash Design layout



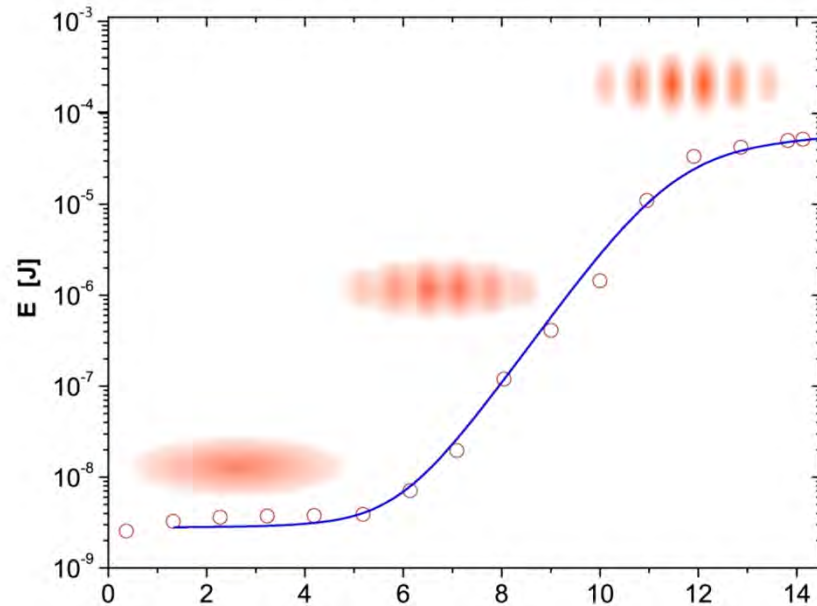
bunch compressor



electron bunch in the undulator



radiation power vs. undulator length



Power grows exponentially with undulator distance z . For a 1-D, mono-energetic beam

$$P \propto e^{\frac{z}{L_{G0}}} \text{ where } L_{G0} = \frac{\lambda_u}{4\sqrt{3\pi\rho}}$$

Pierce Parameter $\rho \sim \frac{\text{radiation power}}{\text{beam power}}$

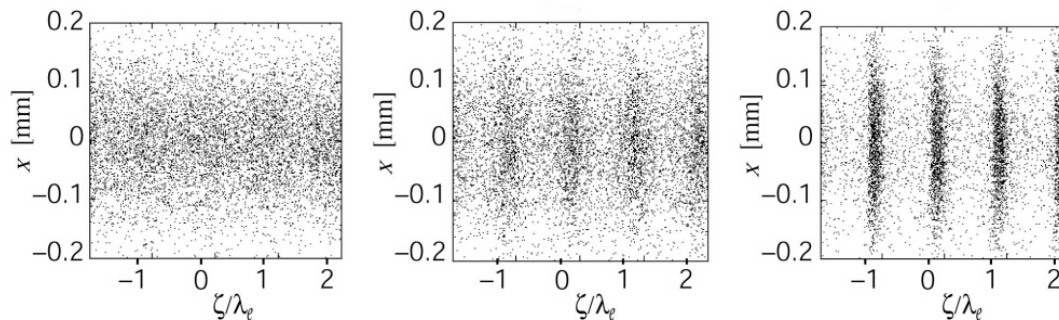
peak current

$$\rho = \frac{1}{4\gamma_0} \left[\frac{I \lambda_u^2 K^2 [JJ]^2}{I_A \pi^2 \epsilon_x \beta_x} \right]$$

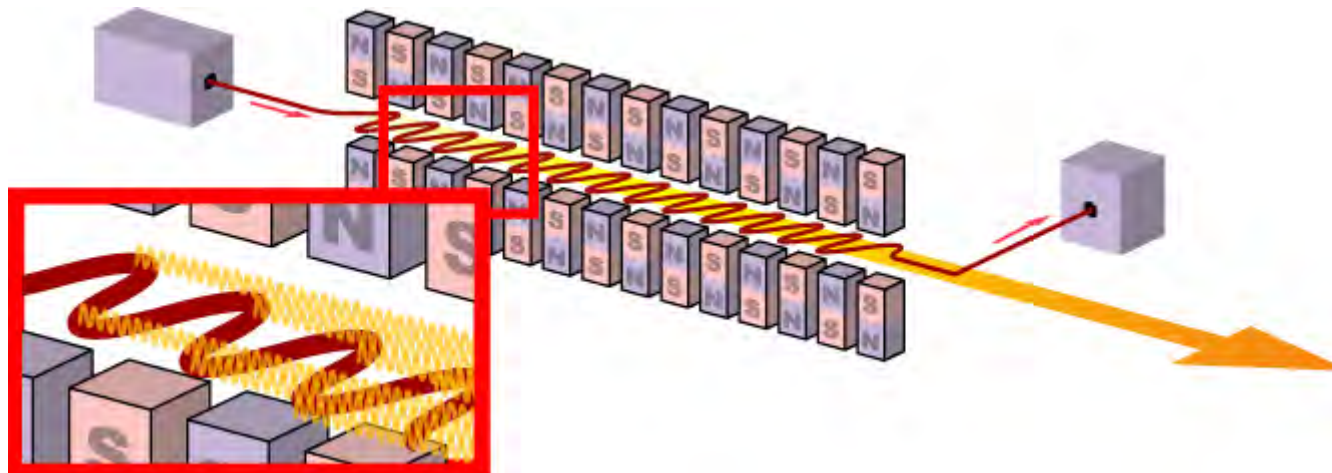
17 kA

emittance

beta function

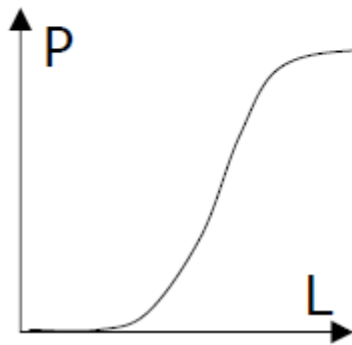
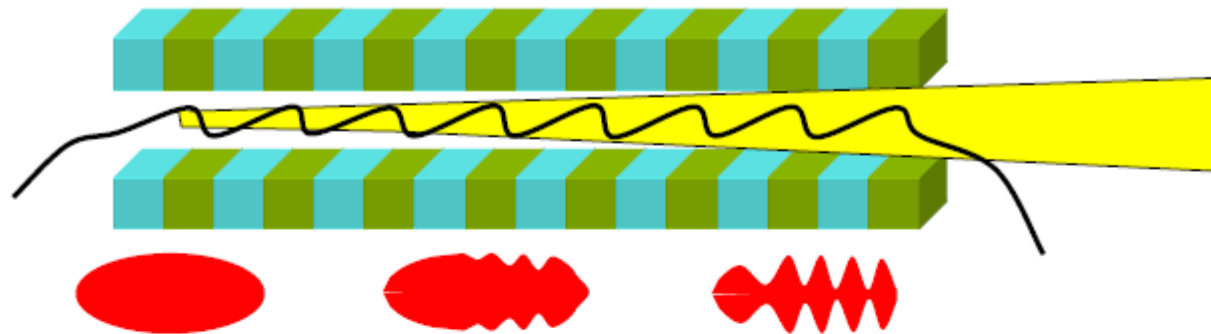


undulator radiation in a long undulator



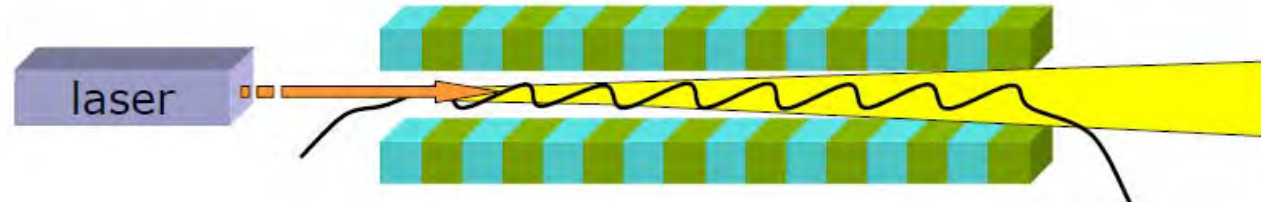
http://en.wikipedia.org/wiki/Free-electron_laser

SASE – Self Amplified Stimulated Emission

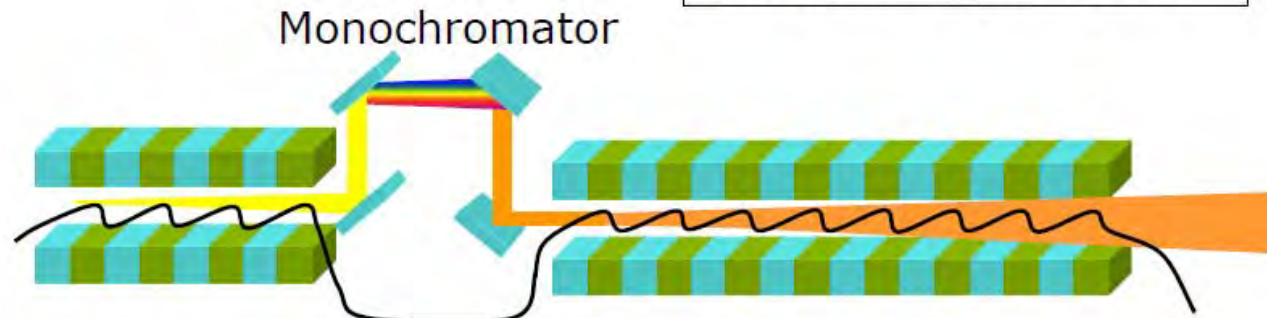


- From noise
- No mirrors (☹ X-rays)
- Tunable
- "Spiky" (t and λ)

Seeding



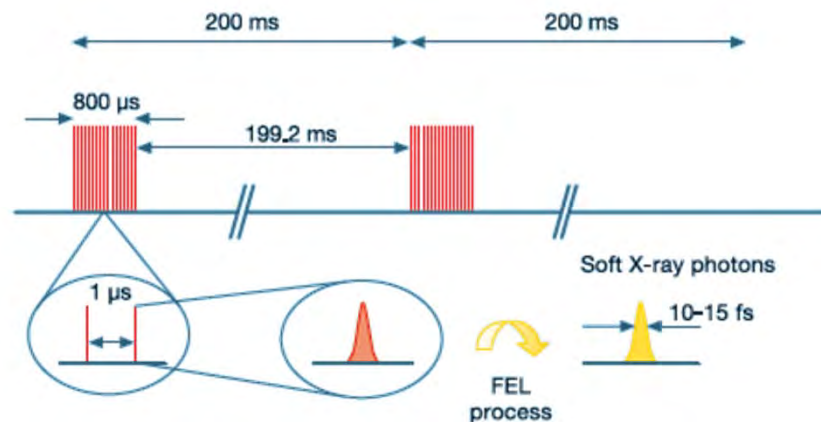
- Remove instabilities
- CHG with gain HGHG



power gain $P(z) = \frac{P_0}{9} \exp(z / L_g) \text{ for } z \geq 2L_g$

Key figures of FLASH

<http://flash.desy.de/>



Performance of the FEL radiation

2005 – 2007

Wavelength range of the fundamental
 Higher harmonics

13 - 47 nm (from fall 2007: 6.5 nm)
 3rd 4.6 nm
 5th 2.7 nm
 (7th 1.9 nm)

Average pulse energy
 Peak pulse energy
 Peak power
 Average power
 Pulse duration
 Spectral width
 Peak brilliance

up to 100 μJ
 170 μJ
 5 GW
 100 mW
 10 - 50 fs
 0.5 - 1%
 $10^{29} - 10^{30}$
 [photons/(s mrad² mm² 0.1% BW)]

http://pr.desy.de/sites2009/site_pr/content/e113/e48/column-objekt221/lbox/infoboxContent222/FLASH_en_2007_ger.pdf

Review of techniques for attosecond X-ray Pulse generation (FEL based)

Main schemes

- emittance spoiler (P. Emma et al. PRL, 2004)
- slicing: wavelength selection (Saldin et al., Opt. Comm. 2004)
- slicing: current enhancement (Zholents et al., PRSTAB 2005)
- slicing: from a Harmonic Cascade FEL (Zholents et al., PRL 2004)
- slicing: energy chirp revisited: tapered undulator (Saldin et al., PRSTAB 2006)
- ultra short electron bunches: single spike (C. Pellegrini et al., 2007)
- energy chirp (C. Schroeder et al., NIMA 2002)
- attosecond trains from mode locking (N. Thompson et al., 2007)

Thank you!

